

# Nonlinear Regge Trajectories in Theory and Practice

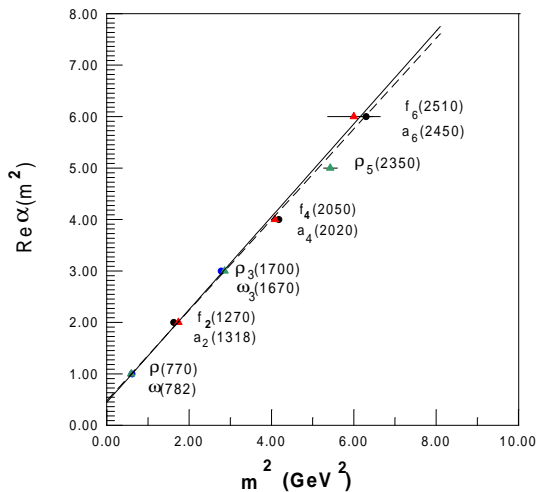
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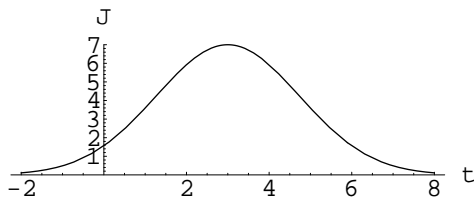
# Spectroscopy



(Desgrolard-Giffon-Martynov-Predazzi, 2000)

# Theory

(pre QCD)  $\lambda\phi^3$ ,  $\lambda\phi^4$  (Lee-Sawyer, 1962)



$$\alpha(t) = \text{const} + \lambda^2 K(t)$$

$$\alpha(t)|_{t \rightarrow -\infty} \rightarrow \begin{cases} -1, & \phi^3 \\ 0, & \phi^4 \end{cases}$$

$$\alpha_P(t) = 1 + \gamma(\sqrt{t_0} - \sqrt{t_0 - t}) \rightarrow -\infty, \quad t \rightarrow -\infty$$

(Fiore-Jenkowszky-Paccanoni-Prokudin, 2003)

$$\alpha(t) = c \ln(b - at), \quad c \sim g^2 \quad (\text{"vector gluon" coupling})$$

(Coon-Suura, 1974)

$$\alpha(t) = a + bt + b\sqrt{(t_0 - t)(t_0^* - t)} \rightarrow \begin{cases} 2bt, & t \rightarrow +\infty \\ -1, & t \rightarrow -\infty \end{cases}$$

(Collins-Kearney, 1984)

+ many others

# QCD

high(-t)

$\alpha_P(t) \rightarrow 1 + \frac{3 \ln 2}{\pi^2} g^2(t) + O(g^{10/3})$

(Kirschner-Lipatov, 1990)

$\alpha_R(t) \rightarrow 0 + O\left(\sqrt{g^2(t)}\right)$

(Kwiecinski, 1982)

low t:

$$\alpha_P(0) = \frac{3 \ln 2}{\pi^2} g^2 \left( 1 - \frac{5}{\pi^2} g^2 \right)$$

(Fadin-Lipatov; Camici-Ciafaloni, 1998)

Renormalization group invariance of a physical quantity P:

$$\left[ \mu^2 \frac{\partial}{\partial \mu^2} + \beta(g^2) \frac{\partial}{\partial g^2} \right] P = 0$$

$$P = \alpha(t) \rightarrow \alpha(t) = \Phi \left[ \frac{t}{\mu^2} e^{K(g^2)} \right]$$

$$\frac{d}{dg^2} K(g^2) = \frac{1}{\beta(g^2)}, \quad \beta(g^2) = -\beta_0 g^4 - \beta_1 g^6 - \dots$$

$$16\pi^2 \beta_0 = 11 - 2n_f/3, \quad 128\pi^4 \beta_1 = 51 - 19n_f/3$$

$$K(g^2) \approx \frac{1}{\beta_0 g^2} + \frac{\beta_1}{\beta_0^2} \ln \frac{1}{g^2} + \dots$$

Simple consequences of RG invariance:

1)  $\alpha(0) = \Phi(0)$ , i.e. the intercept does not depend on  $g^2$

2) If  $\alpha(t)$  is analytic at  $t = 0$  then

$$\alpha(t) = \Phi(0) + \Phi'(0) \frac{t}{\mu^2} e^{K(g^2)} + \dots = \alpha(0) + \alpha'(0)t + \dots$$

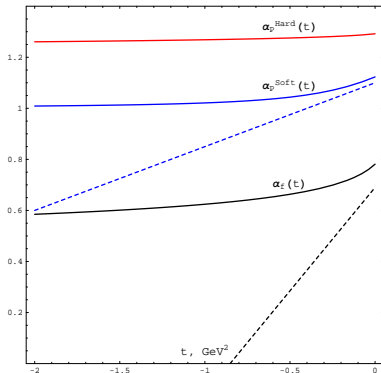
$$\alpha'(0)|_{g^2 \rightarrow 0} \sim \left( \frac{1}{g^2} \right)^{\frac{\beta_1}{\beta_0}} e^{\frac{1}{\beta_0 g^2}}$$

$$\alpha(t) = f[g^2(t)], \quad \alpha(0) = f[g^2(0)] = \Phi \left[ e^{K(g^2(0))} \right]$$

$$K(g^2(0)) = -\infty \quad \left( K(g^2(t)) - K(g^2) = \ln \left( \frac{-t}{\mu^2} \right) \right)$$

$$g^2(0) < \infty \text{ (Shirkov-Solovtsov, 1997)}$$

# Practice

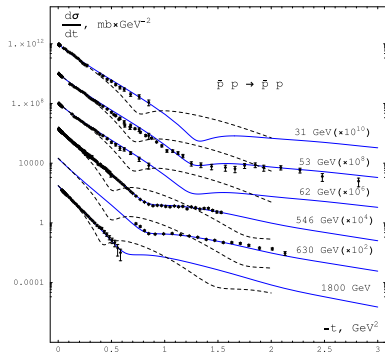
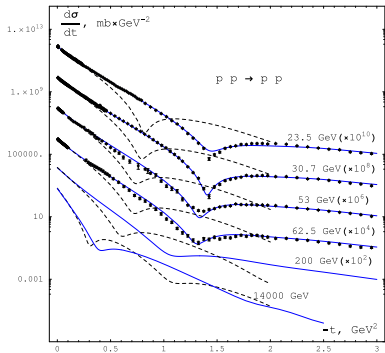


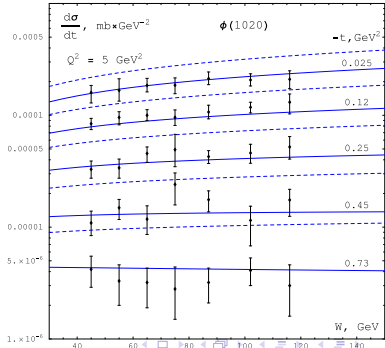
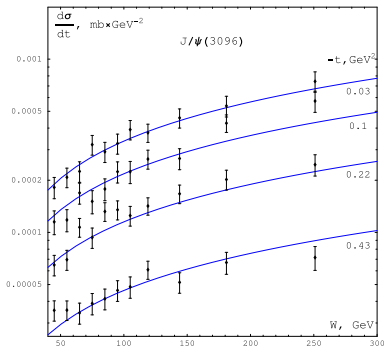
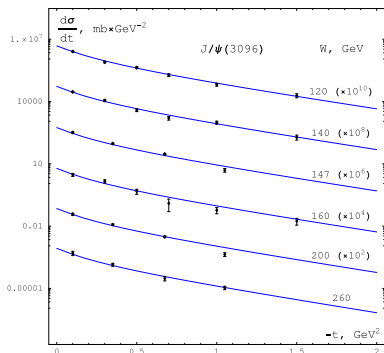
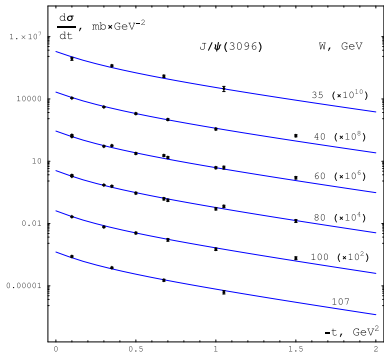
$$\alpha_P^{Soft}(t) = 1 + p_1 \left[ 1 - p_2 t \left( \text{arctg}(p_3 - p_2 t) - \frac{\pi}{2} \right) \right]$$

$$\alpha_P^{Hard}(t) = 1 + \frac{1}{A_H + \left[ \frac{3 \ln 2}{\pi^2} g_s^2(\sqrt{-t + c_H}) \right]^{-1}}$$

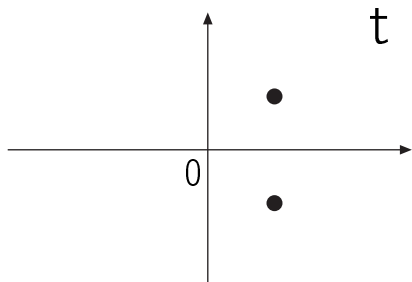
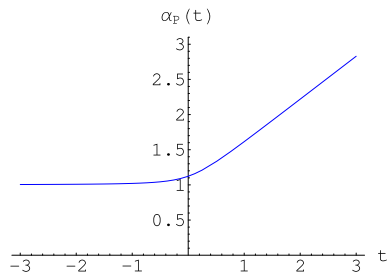
$$\alpha_R(t) = \left( \frac{2}{3\pi^2} g_s^2(\sqrt{-t + c_R}) \right)^{1/2}, \quad g_s^2(\mu) \equiv \frac{(4\pi)^2}{11 - \frac{2}{3}n_f} \left( \frac{1}{\ln \frac{\mu^2}{\Lambda^2}} + \frac{1}{1 - \frac{\mu^2}{\Lambda^2}} \right)$$





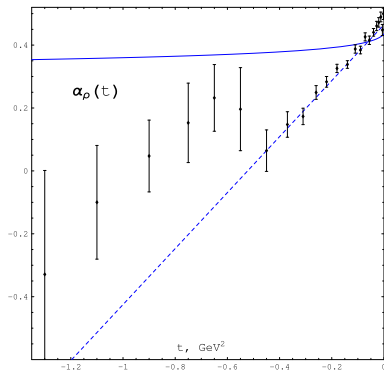


# Problems



Microcausality is under threat? (V.P.-Samokhin, 1990)

# Problems



(Brodsky-Tang-Thorn, 1993; Kaidalov, 2006)

$\pi$ -trajectory?

Heavy quarks?

## (Interim) conclusions:

1. Analyticity of RT in  $t$  seems to imply their singular behavior at  $g^2 \sim 0$ ;
2. QCD (partonlike) asymptotics at high  $(-t)$  and “stringy” asymptotics at high  $t$  imply complex singularities of RT in the  $t$ -plane. Probable clash with microcausality.
3. QCD behavior at high  $(-t)$  does not match with monotony of RT in  $t$  for  $\rho, \pi, \dots$ , heavy quarkonia.